

# Fourier Series

①

Taylor Series  $f(x) = \sum_{n=0}^{\infty} \frac{f^{(n)}(a)}{n!} (x-a)^n$

Fourier Series  $f(x) = \sum_{n=0}^{\infty} \left( a_n \cos \frac{n\pi x}{l} + b_n \sin \frac{n\pi x}{l} \right)$

$f(x)$  is periodic of period  $2l$

let's take a look at the periods

$n$	functions	Periods
0	1	$(2l, 4l, 6l)$
1	$\cos \frac{\pi x}{l}$ $\sin \frac{\pi x}{l}$	$l, (2l, 3l)$
2	$\cos \frac{2\pi x}{l}$ $\sin \frac{2\pi x}{l}$	$\frac{2l}{3}, \frac{4l}{3}, (2l, \frac{8l}{3})$
3	$\cos \frac{3\pi x}{l}$ $\sin \frac{3\pi x}{l}$	

$2l$  is the smallest period shared by all terms.

How to obtain coefficients.

$$a_0 = \frac{1}{2l} \int_{-l}^l f(x) dx$$

$$a_n = \frac{1}{l} \int_{-l}^l f(x) \cos \frac{n\pi x}{l} dx$$

$$b_n = \frac{1}{l} \int_{-l}^l f(x) \sin \frac{n\pi x}{l} dx$$

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Indeed

$$\int_{-l}^l f(x) dx = 2l a_0 + \sum_{n=1}^{\infty} a_n \int_{-l}^l \frac{\cos n\pi x}{l} dx + \sum_{n=1}^{\infty} b_n \int_{-l}^l \frac{\sin n\pi x}{l} dx \quad \rightarrow = 0$$

$$\int_{-l}^l f(x) \frac{\cos n\pi x}{l} dx = a_0 \int_{-l}^l \frac{\cos n\pi x}{l} dx + \sum_{m=1}^{\infty} a_m \int_{-l}^l \frac{\cos m\pi x}{l} \frac{\cos n\pi x}{l} dx + \sum_{m=1}^{\infty} b_m \int_{-l}^l \frac{\cos m\pi x}{l} \frac{\sin n\pi x}{l} dx$$

$$\text{but } \int_{-l}^l \frac{\cos^2 n\pi x}{l} dx = l$$

$$\int_{-l}^l \frac{\sin n\pi x}{l} \frac{\cos n\pi x}{l} dx = 0$$

Same proof for  $b_n$ 

### Fourier Convergence Theorem

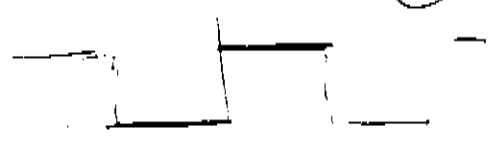
$f$  periodic with period  $2l$ .

$f, f'$  piecewise continuous in  $[-l, l] \Rightarrow$

$\sum_{n=0}^{\infty} \left( a_n \frac{\cos n\pi x}{l} + b_n \frac{\sin n\pi x}{l} \right)$  converges to  $f(x)$  where  $f(x)$  is continuous and to  $\frac{f(x^+) + f(x^-)}{2}$  where  $f$  is discontinuous.

Example

$$f(x) = \begin{cases} 0 & [-l, 0) \\ 4 & [0, l] \end{cases}$$



$$a_0 = \frac{1}{2l} \int_{-l}^l f(x) dx = \frac{1}{2l} \int_{-l}^0 0 dx + \frac{1}{2l} \int_0^l 4 dx = 2$$

↑  
average  
of values  
at  $x=0$

$$a_n = \frac{1}{l} \int_{-l}^l f(x) \cos \frac{n\pi x}{l} dx =$$

$$= \frac{1}{l} \int_{-l}^0 0 dx + \frac{1}{l} \int_0^l 4 \cos \frac{n\pi x}{l} dx = \frac{4}{n\pi} \sin \frac{n\pi x}{l} \Big|_0^l = 0$$

$$b_n = \frac{1}{l} \int_{-l}^l f(x) \sin \frac{n\pi x}{l} dx = \frac{1}{l} \int_0^l 4 \sin \frac{n\pi x}{l} dx =$$

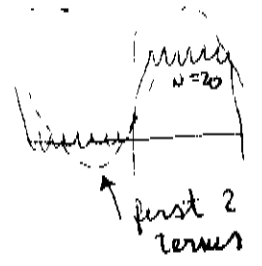
$$= -\frac{4}{n\pi} \left[ \cos \frac{n\pi x}{l} \right]_0^l =$$

$$= \frac{4l}{n\pi} (1 - \cos n\pi)$$

$$b_n = \frac{4l}{n\pi} [1 - (-1)^n] = \begin{cases} \frac{8l}{n\pi} & n=1, 3, \dots \\ 0 & n=2, 4, \dots \end{cases}$$

$$\Rightarrow f(x) = 2 + \frac{8l}{\pi} \sum_{1,3,\dots} \frac{1}{n} \sin \frac{n\pi x}{l}$$

$$= 2 + \frac{8l}{\pi} \sum_1^{\infty} \left( \sin \frac{(2n-1)\pi x}{l} \right) \left( \frac{1}{2n-1} \right)$$



# Gibbs phenomenon

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~~minimum~~ } about 9% overshoot -

What about convergence then? -

Answer: the overshoot point moves  
and convergence is for  $x$  fixed.

Even functions

$$f(x) = f(-x)$$

$$a_0 = \frac{1}{2l} \int_{-l}^l f(x) dx = \frac{1}{l} \int_0^l f(x) dx$$

$$a_n = \frac{2}{l} \int_0^l f(x) \cos \frac{n\pi x}{l} dx$$

$$b_n = \frac{1}{l} \int_{-l}^l f(x) \sin \frac{n\pi x}{l} dx = 0$$

Odd functions.

$$a_0 = 0$$

$$a_n = 0$$

$$b_n = \frac{2}{l} \int_0^l f(x) \sin \frac{n\pi x}{l} dx$$

# Sturm Liouville Theory

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## Sturm Liouville Problem

$$[p(x)y']' + q(x)y + \lambda w(x)y = 0$$

$$\begin{cases} \alpha y(a) + \beta y'(a) = 0 \\ \gamma y(b) + \delta y'(b) = 0 \end{cases} \quad \left\{ \begin{array}{l} w(a)k_a + w(b)k_b = 0 \\ \underline{\underline{w(a)k_a + w(b)k_b = 0}} \end{array} \right.$$

$$\left. \begin{array}{l} p, q, w \text{ continuous} \\ p(x) > 0 \\ w(x) > 0 \end{array} \right\} \text{ on } [a, b]$$

Any  $\lambda$  that allows a non-trivial solution is called an eigenvalue of the problem

$$Ly = \lambda y \quad \cdot \quad L y = \frac{1}{w(x)} [p(x)y']' + \frac{q(x)}{w(x)} y$$

Example:  $y'' + \lambda y = 0$   $y(0) = 0$   
 $y(L) = 0$

$$\Rightarrow y(x) = A \cos \sqrt{\lambda} x + B \sin \sqrt{\lambda} x$$

$\lambda = 0$   $y(x) = A$  ? but for  $\lambda = 0$   $y'' = 0 \Rightarrow y = A + Bx$

$$\begin{array}{l} y(0) = 0 \Rightarrow A = 0 \\ y(L) = 0 \Rightarrow BL = 0 \Rightarrow B = 0 \end{array} \left| \begin{array}{l} \text{solution} \\ \text{is trivial} \end{array} \right.$$

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$$\lambda \neq 0 \quad y(0) = A = 0$$

$$y(L) = 0 = B \sin \sqrt{\lambda} L$$

$B \neq 0$  or we have a trivial solution

$$\Rightarrow \sin \sqrt{\lambda} L = 0 \Rightarrow \sqrt{\lambda} L = n\pi \quad n = 0, \pm 1, \pm 2, \dots$$

Since  $\sin$  is odd use  $n = 0, 1, 2, \dots$

$$\Rightarrow \lambda_n = \frac{n^2 \pi^2}{L^2}$$

↑  
eigenvalue

$$\Rightarrow \phi_n = \sin \frac{n\pi x}{L}$$

↑  
solution  
(eigenfunction)

- 1) Eigenvalues are real
- 2) " " simple (not repeated roots)  
 $\Rightarrow$  one eigenfunction per eigenvalue.

3) Eigenfunctions are orthogonal

4) Eigenfunctions can be used as a basis to obtain an expansion series of  $f$  ( $f$  is piecewise continuous).

$$f(x) = \sum_{n=1}^{\infty} \frac{\langle f, \phi_n \rangle}{\langle \phi_n, \phi_n \rangle} \phi_n(x)$$

$$\langle f, g \rangle = \int_a^b f(x)g(x) \underbrace{w(x)}_{\text{weight function}} dx \quad (7)$$

Periodic BC.

$$\left. \begin{array}{l} y(a) = y(b) \\ y'(a) = y'(b) \end{array} \right\} \text{Does not fall under previous definition.}$$

Singular Cases

$p(x)$  (and/or  $w(x)$ ) vanish at  $a$  or  $b$  (end of interval).

$\Rightarrow$  Need to modify BC.

$$\left. \begin{array}{l} p(a) = 0 \\ p(b) \neq 0 \end{array} \right\} \Rightarrow \begin{array}{l} y \text{ bounded at } a \\ \alpha y(b) + \beta y'(b) = 0 \end{array}$$

$$\left. \begin{array}{l} p(a) = 0 \\ p(b) = 0 \end{array} \right\} \Rightarrow \begin{array}{l} y \text{ bounded at } a \\ y \text{ bounded at } b \end{array}$$

Properties.

- 1) Eigenvalues are real
- 2) Eigenfunctions are orthogonal.

Example

$$y'' + \lambda y = 0 \quad x \in [-L, L]$$

$$y(-L) = y(L) \quad y'(-L) = y'(L)$$

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$$y(x) = \begin{cases} A \cos(\sqrt{\lambda}x) + B \sin(\sqrt{\lambda}x) & \lambda \neq 0 \\ C + Dx & \lambda = 0 \end{cases}$$

$$\lambda = 0 \quad C - DL = C + DL$$

$$\Rightarrow D = 0$$

$$\Rightarrow \boxed{\begin{array}{l} y(x) = C \\ \lambda = 0 \text{ is an} \\ \text{eigenvalue} \end{array}}$$

$\phi_0 = 1$   
eigenfunction.

$$\lambda \neq 0$$

$$\sin(\sqrt{\lambda}L) B = 0$$

$$\sin(\sqrt{\lambda}L) A = 0$$

$$\Rightarrow \sin(\sqrt{\lambda}L) = 0$$

$$\lambda = \left(\frac{n\pi}{L}\right)^2$$

$$\phi_n = \cos \frac{n\pi x}{L}$$

$$\psi_n = \sin \frac{n\pi x}{L}$$

Two  
eigenfunctions



Bessel function example (Singular case) <sup>9</sup>

$$(xy')' + \lambda xy = 0$$

$y(0)$  bounded

$$y(L) = 0$$

$$y(x) = A J_0(\sqrt{\lambda}x) + B Y_0(\sqrt{\lambda}x) \quad \lambda \neq 0$$

$$y(x) = C + D \ln x \quad \lambda = 0$$

~~Case~~

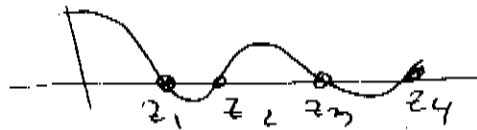
$$\lambda = 0 \quad y(0) \text{ bounded} \Rightarrow D = 0$$

$$y(L) = 0 \Rightarrow C = 0 \Rightarrow \lambda = 0 \text{ not eigenvalue}$$

$$\lambda \neq 0 \quad B = 0 \quad Y_0(0) \text{ is unbounded. } (\rightarrow -\infty)$$

$$y(L) = A J_0(\sqrt{\lambda}L) = 0$$

$\sqrt{\lambda}L$  zero of  $J_0(x)$



$$\lambda_n = \left(\frac{z_n}{L}\right)^2$$

$$\phi_n = J_0\left(z_n \frac{x}{L}\right)$$